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Here is a link towards more detailed lectures delivered during the 2013 VLTI school at Barcelonnette (France): http://hdl.handle.net/2268/155589

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• H. Fizeau and E. Stephan (1868-1870):

"In terms of angular resolution, two small apertures distant of B are equivalent to a single large aperture of diameter B"





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Convolution theorem

$$I(\xi,\eta) = \iint PSF(\xi - \xi', \eta - \eta') O(\xi',\eta') d\xi' d\eta' = PSF(\xi,\eta) \otimes O(\xi,\eta)$$

 $FT(I(\zeta,\eta))(u,v) = FT(PSF(\zeta,\eta))(u,v) \cdot FT(O(\zeta,\eta))(u,v)$

 $\mathbf{u} = \mathbf{B}_{\mathbf{u}} / \lambda, \mathbf{v} = \mathbf{B}_{\mathbf{v}} / \lambda$

 $O(\zeta,\eta) = IFT(FT(O(\zeta,\eta))) = IFT((FT(I(\zeta,\eta)) / FT(PSF(\zeta,\eta))))$

8.2 The convolution theorem

$$f(x) * g(x) = (f * g)(x) = \int_{R^{n}} f(x - t)g(t)dt$$



Convolution product of two 1D rectangle functions. A) f(x), B) g(x), C) g(t) and f(x-t); the dashed area represents the integral of the product of f(x-t) and g(t) for the given x offset, D) $f(x)^*g(x) = (f^*g)(x)$ represents the previous integral as a function of x.



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Camera obscura (cf. shoe box)







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$I(x,y) = Hole(x,y) \otimes Sun(x,y) = \iint Hole(x',y') Sun(x-x',y-y') dx'dy'$



- 1 Introduction
- 2 Reminders
- 3 Brief history of stellar diameter measurements
- 4 Interferometry with two independent telescopes
- 5 Light coherence (Zernicke-van Cittert theorem)
- 6 Examples of optical interferometers
- 7 Results
- 8 Three important theorems (Fundamental theorem, Convolution theorem and Wiener-Khintchin theorem)!

1 Introduction

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1 Introduction

$$\rho = R / z \qquad (1.1)$$



$$F = f / \rho^2$$
 (1.2)

$$T_{eff} = (F/\sigma)^{1/4} = (f/\sigma \rho^2)^{1/4}$$
 (1.3)

Ζ

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- 2 Reminders
- 2.1. Representation of an electromagnetic wave



- 2.1. Representation of an electromagnetic wave
- $E = Re\{ a exp[i2\pi(vt z / \lambda)] \}_{(2.1.3)}$
- E = Re{ a exp[-i ϕ] exp[i $2\pi\nu$ t]} (2.1.4) where $\phi = 2\pi z / \lambda$. (2.1.5)
- $E = a \exp[-i \phi] \exp[i2\pi vt]$

(2.1.6)

An introduction to optical/IR interferometry 2.1. Representation of an electromagnetic wave $E = A \exp[i2\pi vt]$ (2.1.7)with $A = a \exp[-i \phi]$ (2.1.8)

 $v \sim 6 \ 10^{14} \,\text{Hz}$ for $\lambda = 5000 \,\text{\AA}$

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2.1. Representation of an electromagnetic wave

$$\langle E^{2} \rangle = \lim_{T \to \infty} \frac{1}{2T} \int_{-T}^{+T} E^{2} dt \qquad (2.1.9)$$
$$\langle E^{2} \rangle = a^{2} \qquad (2.1.10)$$

 $I = AA^* = |A|^2 = a^2$.

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(2.1.11)

2.2. The Huygens-Fresnel principle



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2.2. The Huygens-Fresnel principle







2.2. The Huygens-Fresnel principle



3 Brief history of stellar diameter measurements



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3 Brief history of stellar diameter measurements

b) Newton:

$$V_{\odot} - V = -5 \log (z / z_{\odot}), \quad (3.1)$$
$$\Delta = 2 R_{\odot} / z \qquad , \quad (3.2)$$

$$\Lambda \sim 2.10^{-3'}$$
 (8 10^{-3''}) (3.3)

c) Fizeau-type interferometry

4 Interferometry with two independent telescopes



4 Interferometry with two independent telescopes

b) Fizeau ... the father of stellar interferometry (1868) If $\Delta \ge \phi/2 = \lambda / (2B)$, (4.7) fringe disappearance!

Fringe visibility:

$$\upsilon = \left(\frac{I_{\max} - I_{\min}}{I_{\max} + I_{\min}}\right)$$



4 Interferometry with two independent telescopes

b) Fizeau ... the father of stellar interferometry (1868)



4 Interferometry with two independent telescopes

b) Fizeau ... the father of stellar interferometry (1868)



Stéphan, 1873 Δ << 0.16''

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Marseille 80 cm telescope



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4 Interferometry with two independent telescopes
 b) Fizeau ... the father of stellar interferometry (1868)



• Michelson and Pease (1920)



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4 Interferometry with two independent telescopes
 b) Fizeau ... the father of stellar interferometry (1868)

- Anderson
- Brown and Twiss (1956)
- Radio Interferometry (1950)



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- 5 Light coherence
- 5.1 Quasi monochromatic light (waves and wave groups)



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- 5 Light coherence
- 5.1 Quasi monochromatic light (waves and wave groups)



5 Light coherence

5.1 Quasi monochromatic light (waves and wave groups)



- 5 Light coherence
- 5.1 Quasi monochromatic light (waves and wave groups)

 $λ_0 = 2.2μm$ λ∈ [2.07 ; 2.33]μm Δλ = 0.13μm









- 5 Light coherence
- 5.2 Fringe visibility



- 5 Light coherence
- 5.2 Fringe visibility

$$I_{q} = I + I + 2 I \operatorname{Re} \{ \gamma_{12}(\tau) \} \qquad \gamma_{12}(\tau) = \langle V_{1}^{*}(t) V_{2}(t-\tau) \rangle$$
(5.2.5)
(5.2)

$$\gamma_{12}(\tau) = \langle A_1^*(z,t) A_2(z,t-\tau) \rangle \exp(-i2\Pi v\tau) / I$$
(5.2.7)



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 $\gamma_{12}(\tau) = |\gamma_{12}(\tau=0)| \exp(i\beta_{12} - i2\Pi v\tau)$

.6)

(5.2.8)

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- 5 Light coherence
- 5.2 Fringe visibility

$$I_{q} = I + I + 2I |\gamma_{12}(0)| \cos(\beta_{12} - 2\Pi \nu \tau)$$
(5.2.9)

$$\upsilon = \left(\frac{I_{\text{max}} - I_{\text{min}}}{I_{\text{max}} + I_{\text{min}}}\right) = |\gamma_{12}(0)|$$

(5.2.10)

An introduction to optical/IR interferometry Brief summary of some main results:

 $\rho = R / z$



 $T_{eff} = (F/\sigma)^{1/4} = (f / \sigma \rho^2)^{1/4}$

 $E = A(z) \exp[i2\pi vt]$

 $E = A(z, t) \exp[i2\pi v t]$

$$\tau = 1 / \Delta v$$
 $\lambda_{\rm eff} = \lambda^2 / \Delta \lambda$

I = A A* = $|A|^2$ = a^2 . 2017 Evry Schatzman School - Roscoff

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If $\Delta \ge \lambda / (2B)$, fringe disappearance!



$$I_{q} = I + I + 2I |\gamma_{12}(0)| \cos(\beta_{12} - 2\Pi v\tau)$$

$$\gamma_{12}(\tau) = \langle V_1^*(t) V_2(t-\tau) \rangle / I$$

Fringe visibility: $v = \left(\frac{I_{\text{max}} - I_{\text{min}}}{I_{\text{max}} + I_{\text{min}}}\right) = |\gamma_{12}(0)|$

- 5 Light coherence
- 5.3 Spatial light coherence

??
$$\gamma_{12}(\tau = 0) = \langle V_1^*(t) V_2(t) \rangle / I_{(5.3.1)}$$

$$V_{1}(t) = \sum_{i=1}^{N} V_{i1}(t)$$

$$V_{2}(t) = \sum_{i=1}^{N} V_{i2}(t)$$

$$I_{q}?$$

$$S = \Sigma dS_{i}$$
for i = 1, N
$$P_{1}$$

$$V_{1}(t)$$

$$V_{2}(t) = \sum_{i=1}^{N} V_{i2}(t)$$

$$P_{2}$$

$$V_{2}(t)$$

$$V_{2}(t)$$

$$V_{2}(t)$$

$$V_{2}(t)$$

$$V_{12}(0) = \left[\sum_{i=1}^{N} \langle V_{1}^{*}V_{2} \rangle + \sum_{i=1}^{N} \langle V_{1}V_{1} \rangle \langle V_{1} \rangle \right] = \left[\sum_{i=1}^{N} \langle V_{1}^{*}V_{2} \rangle + \sum_{i=1}^{N} \langle V_{1}V_{1} \rangle \langle V_{1} \rangle \langle V_{1} \rangle \right] = \left[\sum_{i=1}^{N} \langle V_{1}^{*}V_{2} \rangle + \sum_{i=1}^{N} \langle V_{1}V_{1} \rangle \langle V_{1} \rangle \langle V_{1} \rangle \langle V_{1} \rangle \right] = \left[\sum_{i=1}^{N} \langle V_{1}V_{1} \rangle \langle V_{1} \rangle \langle V_{1}$$

5 Light coherence

5.3 Spatial light coherence

 $\gamma_{12}(0) = \left[\sum_{i=1}^{N} \langle V_{i1}^* V_{i2} \rangle + \sum_{i=1}^{N} \langle V_{i1} V_{j2} \rangle\right] / I$

$$V_{i1}(t) = \left(\frac{\alpha_i(t - \gamma_{i1}/c)}{\gamma_{i1}}\right) \exp\left\{i2\Pi v(t - \gamma_{i1}/c)\right\}_{(5.3.4)}$$
$$V_{i2}(t) = \left(\frac{\alpha_i(t - \gamma_{i2}/c)}{\gamma_{i2}}\right) \exp\left\{i2\Pi v(t - \gamma_{i2}/c)\right\}$$

$$V_{i1}^{*}(t)V_{i2}(t) = \left| \frac{\alpha_{i}(t - r_{i1}/c)}{(r_{i1}r_{i2})} \exp\left\{-i2\Pi v(r_{i2} - r_{i1})/c\right\} (5.3.5)$$

as long as:
$$|r_{i1} - r_{i2}| \le c/\Delta v = \lambda^{2}/\Delta \lambda = \ell$$
(5.3.6)

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- 5 Light coherence
- 5.3 Spatial light coherence

$$I(s)ds = \left|a_i(t-r/c)\right|^2$$

$$\gamma_{12}(0) = \int_{S} \frac{I(s)}{r_{1}r_{2}} \exp\{-i2\Pi(r_{2} - r_{1})/\lambda\} ds / I$$
(5.3.8)

!!! Theorem of Zernicke-van Cittert !!!

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(5.3.7)

- 5 Light coherence
- 5.3 Spatial light coherence



 $|\mathbf{r}_2 - \mathbf{r}_1| = |\mathbf{P}_2 \mathbf{P}_i - \mathbf{P}_1 \mathbf{P}_i| = |-(\mathbf{X}^2 + \mathbf{Y}^2) / 2 \mathbf{Z'} + (\mathbf{X} \boldsymbol{\zeta} + \mathbf{Y} \boldsymbol{\eta})|$ (5.3.9)

where $\zeta = X' / Z'$ and $\eta = Y' / Z'$

(5.3.10)

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- 5 Light coherence
- 5.3 Spatial light coherence

$$\gamma_{12}(0, X/\lambda, Y/\lambda) = \exp(-i\phi_{X,Y}) \frac{\iint I(\xi, \eta) \exp\{-i2\Pi(X\xi + Y\eta)/\lambda\} d\xi d\eta}{\iint I(\xi', \eta') d\xi' d\eta'}$$
(5.3.11)

$$I'(\zeta,\eta) = I(\zeta,\eta) / \iint_{S} I(\zeta',\eta') d\zeta' d\eta'$$
(5.3.12)

Setting $u = X/\lambda$, $v = Y/\lambda$:

$$\gamma_{12}(0,u,v) = \exp(-i\phi_{u,v}) \iint_{S} I'(\zeta,\eta) \exp\{-i2\Pi(u\zeta+v\eta)\} d\zeta d\eta$$
(5.3.13)

$$I'(\xi,\eta) = \iint \gamma_{12}(0,u,v) \exp\{i \phi_{u,v}\} \exp\{i 2\Pi(\xi u + \eta v)\} d(u) d(v)$$
(5.3.14)

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5.4 Fourier transform (cf. Léna 1996) 5.4.1 Definitions:

$$FT_f(s) = \int_{-\infty}^{\infty} f(x) e^{-2i\pi sx} dx,$$

$$f(x) = \int_{-\infty}^{\infty} FT f(s) e^{2i\pi sx} ds$$

$$\int_{-\infty}^{\infty} \left| f(x) \right|^2 dx.$$

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(5.4.1)

(5.4.2)

(5.4.3)

5.4 Fourier transform (cf. Léna 1996)
 5.4.1 Definitions: Generalisation:

$$FT_f(\vec{w}) = \int_{-\infty}^{\infty} f(\vec{r}) e^{-2i\pi \vec{r} \cdot \vec{w}} d\vec{r} \quad (5.4.4)$$

5.4.2 Some properties: a) Linearity: $FT_(af) = a FT_f$, $a \in \Re$, a being a constant, (5.4.5)

(5.4.6)

5.4 Fourier transform (cf. Léna 1996)
 5.4.2 Some properties: b) Symmetry & parity:



5.4 Fourier transform (cf. Léna 1996)
c) Similarity:

$$FT_{(f(x/a))(s)} = |a| FT_{(f(x))(sa)},$$

(5.4.9)

where $a \in \Re$, is a constant.

d) Translation: $FT_{(f(x - a))(s)} = e^{-2i\pi as} FT_{(f(x))(s)}$

(5.4.10)

5.4 Fourier transform (cf. Léna 1996)

e) Derivation:

- $\begin{array}{l} {\sf FT}_{df/dx}(s) = 2i\pi s \; {\sf FT}_{f}(s), \; {\sf FT}_{d^nf/dx^n}(s) = (2i\pi s)^n \\ {\sf FT}_{f}(s). \end{array} \tag{5.4.11}$
- 5.4.3 Some important cases (one dimension): a) Door (top-hat) function: $\Pi(x) = 1$ if $x \in]-1/2$, 1/2[, (5.4.12) = 0 if $x \in]-\infty$, -1/2] or $x \in [1/2, \infty[$.

5.4 Fourier transform (cf. Léna 1996)

The door function and its Fourier transform (cardinal sine) -



FT_ ($\Pi(\mathbf{x})$)(s) = sinc(s) = sin(π s) / π s. (5.4.13)

 $FT_{(\pi(x/a))(s)} = |a| sinc(as) = |a| sin(\pi as) / \pi as.$ (5.4.14)

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5.4 Fourier transform (cf. Léna 1996)
 b) Dirac distribution:

$$\delta(x) = \int_{-\infty}^{\infty} e^{2i\pi sx} ds$$

its Fourier transform is thus unity (= 1) in the interval]- ∞ , ∞ [.

(5.4.15)

- 5 Light coherence
- 5.5 Aperture synthesis

$$\upsilon = \left| \gamma_{12}(0, B/\lambda) \right| = \left| \frac{\sin(\Pi Bb/\lambda z')}{\Pi Bb/\lambda z'} \right|$$
(5.5.2)

 $\Pi Bb / \lambda z' = \Pi (5.5.3)$

 $\Delta \sim \lambda / B$, for a (5.5.4) rectangular source.

 $\Delta \sim 1.22 \lambda / B$, for^(5.5.5) a circular source !



- 5 Light coherence
- 5.5 Aperture synthesis

Exercises (...): point-like source?, double pointlike source with a flux ratio = 1?, gaussian-like source?, uniform disk source?, ...

Case of a double point-like source with a flux ratio = 1



Case of a double point-like source with a flux ratio 0.7/0.3

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Variation of the fringe contrast as a function of the angular separation between the two stars:



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$$\upsilon = \left(\frac{I_{\text{max}} - I_{\text{min}}}{I_{\text{max}} + I_{\text{min}}}\right) = \left|\gamma_{12}(0)\right| = TF(I) = \frac{2J_1(\pi\theta_{UD}B/\lambda)}{\pi\theta_{UD}B/\lambda}$$





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For the case of the Sun:

 $\vartheta_{UD} = 1.22\lambda/B = 1.22\ 0.55/B(\mu) = 30' \ x \ 60''/206265$ B(\mu) = 206265 x 1.22 x 0.55/(30 x 60) = 76.9 \mu d(\mu) = 7.2 or 14.4 \mu $\Rightarrow \sigma = 2.44 \ \lambda/d = 7.8^{\circ} \text{ or } 3.9^{\circ}$

See the masks!


First fringes on the Sun: 9/4/2010

 $B = 29.4 \mu$ d = 11.8 μ





OVLA_Sun_2











Interferometric observations on 10/4/2010 of Procyon, Mars and Saturn, using the 80cm telescope at Haute-Provence Observatory and adequate masks (coll. with Hervé le Coroller) ... 25-29/9/2017 2017 E







Procyon B = 12 mm d = 2 mm



Mars B = 12 mmd = 2 mm



Saturn B = 12 mmd = 2 mm



An introduction to optical/IR interferometry Brief summary of main results:



$$V = \left| \gamma_{12}(0, u, v) \right| = \left| \iint_{S} I'(\zeta, \eta) \exp\left\{ -i2\Pi \left(u\zeta + v\eta \right) \right\} d\zeta d\eta$$

 $I'(\zeta,\eta) = \iint \gamma_{12}(0,u,v) \exp\left\{i2\Pi(\zeta u + \eta v)\right\} d(u)d(v)$

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6 Some examples of optical interferometers





First fringes with I2T









7 Some results

Star	Spectral type	Luminosity class	Angular diameter × 10 ⁻³ seconds of arc 20		
α Boo	K2	Giant			
a Tau	K5	Giant	20		
a Sco	M1-M2	Super-giant	40		
8 Peg	M2	Giant	21		
o Cet	M6e	Giant	47		
α Ori	M1-M2	Super-giant variable	34-→47		

7 Some results

NOM SI		DIAMÈTRE	MESURÉ	R/R ©	TEMPÉRATURE EFFECTIVE		DISTANCE
	SPECTRE A	λ == 0,55 μm en ma. d'erc			λ = 0.55 μm en degrés Kelvin). – 2,2 µm en degrés Kalvin	en persecs (1 pc - 3,25 el
a Cas	кон	54±0.6		24 ± č	4700 ± 300	24.2. 42	45±9
8 And	MOIII	13.2 ± 1.7	14.4 ± 0.5	33±9	3900 ± 260	3711±64	23±3
y And	K38	6.8±0.5		50±14	4650 ± 250	1012-01-020100	西土场
a Per	FSIb	2.9±0.4		\$5 主9	7000±500		176±6
a Cyg	A2la	27±03		145±45	\$200 ± 600	{	500±100
a Ari	K260	7, \$ ±1		15±5	4300±350		23 ± 4
8 Gem	KONI	7.8±0.8		#±2	4900 ± 220		j 14±1
8 Umi	×411	\$.\$±1		30±9	4229±300		31±11
y Ora	×581	8.7±0.8	10.2±1.4	45 :: 10	4300±230	3960±270	59±21
5 Dra	6911	3.8±0.3	i	16±5	4530±220		36 ± 8
u Gera	MSBI		14.6 ± 0.8	94±30		3860±95	60±15
a Teu	×5ŧII	-	20.7±0.4	47 ± 7		3904±34	21±3
a Boo	K2111		21.5±1.2	25±6		4240±120	\$1±2
a Aur.	GSIII	8.0±1.2	[11.7±2	5400 ± 200		137±06
a Aur.	GOHI	4.8 ± 1.5		7.1±2	5950 ± 200	F	13.7±0.6
alvr	AQV	3.0±0.2		2.6±0.2			8.1 主 0.3

6 Some examples of optical interferometers





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http://www.aeos.uig.ac.6ef/HARI/

 6 Some examples of optical interferometers Interferometry to-day is:

Very Large Telescope Interferometer (VLTI)

- 4 x 8.2m UTs
- 4 x 1.8m ATs
- Max. Base: 200m









6 Some examples of optical interferometers



 6 Some examples of optical interferometers Interferometry to-day is also:

The CHARA interferometer

6 x 1m
telescopes
Max. Base:
330m





 6 Some examples of optical interferometers Interferometry to-day is also:

Palomar Testbed Interferometer (PTI)

3 x 40cm
telescopes
Max. Base:
110m

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6 Some examples of optical interferometers

Interferometry to-day

is also:

Keck interferometer

2 x 10m
telescopes
Base: 85m





6 Other examples of interferometers: ALMA



6 Other examples of interferometers: DARWIN



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8 Three important theorems ... and some applications

8.1 The fundamental theorem

8.2 The convolution theorem

8.3 The Wiener-Khintchin theorem

Réf.: P. Léna; Astrophysique: méthodes physiques de l'observation (Savoirs Actuels / CNRS Editions)

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8.1 The fundamental theorem

 $a(p,q) = TF_(A(x,y))(p,q),$ $a(p,q) = \int_{R^2} A(x,y) \exp[-i2\pi(px+qy)] dx dy,$

with

 $p = x' / (\lambda f)$ $q = y' / (\lambda f)$



8.1 The fundamental theorem

The distribution of the complex amplitude a(p,q) in the focal plane is given by the Fourier transform of the distribution of the complex amplitude A(x,y) in the entrance pupil plane.

8.1 The fundamental theorem Application: Point Spread Function determination



$$A(x,y) = A_0 P_0(x,y),$$
 (8.1.1)

$$P_0(x,y) = \Pi(x / a) \Pi(y/a).$$
 (8.1.2)

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8.1 The fundamental theorem $a(p, q) = TE \left[A(x, y)\right](p, q) = \int \int A \exp[-i2\pi(px + qy)](p, q) dx$

$$a(p,q) = TF [A(x,y)](p,q) = \int_{-a/2} \int_{-a/2} A_0 \exp[-i2\pi(px+qy)]dxdy$$
(8.1.3)

$$a(p,q) = A_0 \int_{-a/2}^{a/2} \exp\left[-i2\pi px\right] dx \int_{-a/2}^{a/2} \exp\left[-i2\pi qy\right] dy$$
(8.1.4)

 $a(p,q) = A_0 a^2 [sin(\pi pa) / (\pi pa)] [sin(\pi qa) / (\pi qa)].$ (8.1.5)

 $i(p,q) = a(p,q) a^{*}(p,q) = |a(p,q)|^{2} = |h(p,q)|^{2} =$ $= i_{0} a^{4} [sin(\pi pa) / (\pi pa)]^{2} [sin(\pi qa) / (\pi qa)]^{2}.$ (8.1.6)
8.1 The fundamental theorem Application: Point Spread Function determination



 $\Delta p = \Delta x' / (\lambda f); \Delta q = \Delta y' / (\lambda f) = 2/a \Rightarrow \Delta \phi_{x'} = \Delta \phi_{y'} = 2\lambda/a \quad (8.1.7)$ 25-29/9/2017 2017 Evry Schatzman School - Roscoff



BESSEL FUNCTIONS (REMINDER)

Integral representation of the Bessel functions

$$J_0(x) = \frac{1}{\pi} \int_0^{\pi} \cos[x\sin(\vartheta)] d\vartheta \qquad \qquad J_n(x) = \frac{1}{\pi} \int_0^{\pi} \cos[n\vartheta - x\sin(\vartheta)] d\vartheta$$

Undefined integral

$$\int x' J_0(x') dx' = x J_1(x)$$

Graphs of the $J_0(x)$ and $J_1(x)$ functions



Series development (x ~ 0): $J_0(x) = 1 - x^2/2^2 + x^4/(2^24^2) - x^6/(2^24^26^2) + \dots$ $J_1(x) = x/2 - x^3/(2^24) + x^5/(2^24^26) - x^7/(2^24^26^28) + -1$ $J_n(x) = (2 / (\pi x))^{1/2} \cos(x - n\pi/2 - \pi/4) \dots$ and when x is large!

8.1 The fundamental theorem: 2 telescope interferometer



Two coupled optical telescopes: simplified optical scheme (a). Distribution of the complex amplitude for the case of two circular (b) op square (c) apertures and 127 corresponding impulse response (d).

8.1 The fundamental theorem: 2 telescope interferometer

$$h(p,q) = TF(P(x,y)(p,q) = \int_{R^2} P(x,y) \exp[-i2\pi(px+qy)] dx dy \quad (8.1.10)$$

 $h(p,q) = TF(P_0(x+D/2) + P_0(x-D/2))(p,q) =$ $TF(P_0(x+D/2))(p,q) + TF(P_0(x-D/2))(p,q) =$ $\exp(i\pi D) TF(P_0(x))(p,q) + \exp(-i\pi D) TF(P_0(x))(p,q) =$ $(\exp(i\pi D) + \exp(-i\pi D)) TF(P_0(x))(p,q) =$ $2\cos(\pi D) TF(P_0(x))(p,q)$

(8.1.11)

For the particular case of two square apertures: $i(p,q) = \left| h(p,q) \right|^{2} = 4\cos^{2}(\pi pD) d^{4} \left(\frac{\sin(\pi qd)}{\pi qd} \right)^{2} \left(\frac{\sin(\pi pd)}{\pi pd} \right)^{2} \left(\frac{\sin(\pi pd)}{\pi pd} \right)^{2}$ 25-29/9/2017

(8.1.12)



Delay lines at the VLTI





How are those locations related to the *uv* coverage?



Examples of Fourier plane coverage

Dec -15

Dec -65



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How does the uv plane coverage impact imaging?



8.2 The convolution theorem

 $e(p,q) = O(p,q) * |h(p,q)|^2$,



$$e(p,q) = \int_{\mathbb{R}^2} O(r,s) \left| h(p-r,q-s) \right|^2 drds$$

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8.2 The convolution theorem

For the case of a point-like source:

 $O(p,q) = E \delta(p,q),$

(8.2.1)

 $\delta(x) = 0$ if $x \neq 0$, $\delta(x) = \infty$ if x = 0] and

(8.2.2)

 $e(p,q) = O(p,q) * |h(p,q)|^2 = E \delta(p,q) * |h(p,q)|^2 = E |h(p,q)|^2 (8.2.3)$

8.2 The convolution theorem

More generally, since FT_(f * g) = FT_(f) FT_(g).

We find, because $e(\zeta,\eta) = O(\zeta,\eta) * PSF(\zeta,\eta)$ (8.2.4)

(8.2.5)

that: FT_(e(ζ,η)) = FT_(O(ζ,η)) FT_(PSF(ζ,η)), (8.2.6)

and, finally, $O(\zeta,\eta) = FT^{-1}[FT_(e(\zeta,\eta)) / FT_PSF(\zeta,\eta)].$ (8.2.7)

8.2 The convolution theorem

$$O(p,q) = (\lambda^{2} E / \phi^{2}) \Pi(p \lambda / \phi) \Pi(q \lambda / \phi).$$

$$e(p,q) = O(p,q) * |h_{0}(p,q)|^{2}.$$

$$e(p) = O(p) * |h_{0}(p)|^{2},$$

$$e(p) = 2d^{2}(\lambda/\phi)\sqrt{E} \int_{p-\phi/2\lambda}^{p+\phi/2\lambda} \left(\frac{\sin(\pi r d)}{\pi r d}\right)^{2} \cos^{2}(\pi r D) dr$$
(8.2.10)

$$\begin{pmatrix} \sin(\pi rd) \\ \pi rd \end{pmatrix} \approx \text{Cte sur } [p-\phi/2\lambda, p+\phi/2\lambda], \quad \text{et} \quad (8.2.11)$$

$$e(p) = 2 d^{2}(\lambda/\phi) \sqrt{E} \left(\frac{\sin(\pi pd)}{\pi pd} \right)^{2} \int_{p-\phi/2\lambda}^{p+\phi/2\lambda} \cos^{2}(\pi rD) dr. \quad (8.2.12)$$

$$25-29/9/2017 \qquad (8.2.12)$$

8.2 The convolution theorem

$$e(p) = 2d^{2} \left(\frac{\sin(\pi pd)}{\pi pd}\right)^{2} [O(p) * \cos^{2}(\pi pD)], \qquad (8.2.13)$$

$$e(p) = 2d^{2} \left(\frac{\sin(\pi pd)}{\pi pd}\right)^{2} \left[\frac{1}{2}\int_{R} O(p)dp + \frac{1}{2}O(p) * \cos(2\pi pD)\right]$$
(8.2.14)

$$e(p) = A \left[B + \frac{1}{2} \operatorname{Re}(O(p) * \exp(i2\pi pD)) \right],$$
 (8.2.15)

$$A = 2d^{2} \left(\frac{\sin(\pi pd)}{\pi pd}\right)^{2} et \quad B = \frac{1}{2} \int_{R} O(p) dp,$$
25-29/9/2017 et $B = \frac{1}{2} \int_{R} O(p) dp,$
2017 Evry Schatzman School - Roscoff (8.2.16)

8.2 The convolution theorem

$$e(p) = A \left[B + \frac{1}{2} \operatorname{Re} \left(\int_{R} O(r) \exp(i2\pi(p-r)D) dr \right) \right], \qquad (8.2.17)$$

$$e(p) = A \left| B + \frac{1}{2} \cos(2\pi pD) TF (O(r))(D) \right|,$$
 (8.2.18)

 $\gamma(D) = (e_{\text{max}} - e_{\text{min}}) / (e_{\text{max}} + e_{\text{min}}),$

(8.2.19)

 $\gamma(D) = TF_(O(r))(D) / (2B) = TF_(O(r))(D) / \int O(p) dp.$ (8.2.20)

8.3 The Wiener-Khintchin theorem
In our case, this theorem merely states that the Fourier tranform of the PSF (see Eq. (8.2.7)) is the auto-correlation function of the distribution of the complex amplitude in the pupil plane:

$$TF(|h(p,q)|^{2}) = \iint A^{*}(x,y) A(x+p,y+q) dx dy$$



Diagram representing the autocorrelation as a function of the space frequency, for a 2 telescope interferometer, each having a diameter a, separated by a baseline b. The autocorrelation of ^{25-29/9/2017} pupil give access²⁰¹⁷ Fryschatzman School Rescoff enders.¹⁴⁸